

光電効果

束縛状態にある電子が光を吸収し、散乱状態に移る過程

ここで、静止エネルギー ⇒ 飛び出た電子の運動エネルギー ⇒ 束縛エネルギー

つまり、 $m c^2 \gg k \gg l$ を仮定する。

$$\begin{aligned} H &= \frac{(P - eA)^2}{2m} - \frac{e^2}{r} - \frac{e}{2m} \nabla \cdot H \\ &= \frac{P^2}{2m} - \frac{Z e^2}{r} - \frac{e}{m} P \cdot A + \frac{e^2}{2m} A^2 - \frac{e}{2m} \nabla \cdot H \quad \text{≈ あてたい} \end{aligned}$$

スピノル無視すると、 $H_{\text{int}} = -\frac{e}{m} P \cdot A$

電磁場の量子化は

$$\begin{aligned} A &= \sum_{k, \lambda} \frac{1}{\sqrt{2\omega_k V}} \epsilon(k, \lambda) (c_{k, \lambda} e^{-ikx} + c_{k, \lambda}^+ e^{ikx}) \\ &= \sum_{k, \lambda} \frac{1}{\sqrt{2\omega_k V}} \epsilon(k, \lambda) (c_{k, \lambda} e^{ikx} e^{-ik\omega_k t} + c_{k, \lambda}^+ e^{-ikx} e^{ik\omega_k t}) \end{aligned}$$

↑ イネルギー保存の一項に加わる。

始状態は $\langle f | i \rangle = \frac{1}{\sqrt{\pi}} \left(\frac{z}{a_0}\right)^{3/2} e^{-\frac{z}{a_0} r} | 0 \rangle$, 終状態は $\langle x | f \rangle = \frac{1}{\sqrt{V}} e^{iP_f x} | 0 \rangle$

よって $\langle f | H_{\text{int}} | i \rangle = -\frac{e}{m} \langle f | A \cdot P | i \rangle$

$$\begin{aligned} &= -\frac{e}{m} \int d^3x \int d^3x' \langle f | x \rangle \langle x | A \cdot P | x' \rangle \langle x' | i \rangle \\ &= -\frac{e}{m} \int d^3x \int d^3x' \frac{e^{-iP_f x}}{\sqrt{V}} \langle 0 | \left(\sum_{k, \lambda} \frac{1}{\sqrt{2\omega_k V}} \epsilon(k, \lambda) c_{k, \lambda} e^{ikx'} \right) \cdot (-i\nabla') \left[\frac{1}{\sqrt{\pi}} \left(\frac{z}{a_0}\right)^{3/2} e^{-\frac{z}{a_0} r} \right] \\ &\quad \times \delta^{(3)}(x - x') c_{k, \lambda}^+ | 0 \rangle \end{aligned}$$

$$= -\frac{e}{m} \int d^3x \frac{e^{-iP_f x}}{\sqrt{V}} \sum_{k, \lambda} \delta_{kk} \delta_{\lambda \lambda} \frac{1}{\sqrt{2\omega_k V}} e^{ikx} \epsilon(k, \lambda) \cdot (-i\nabla') \left[\frac{1}{\sqrt{\pi}} \left(\frac{z}{a_0}\right)^{3/2} e^{-\frac{z}{a_0} r} \right]$$

$$= -\frac{e}{m} \cdot \frac{-i}{\sqrt{2\omega_k V^2 \pi}} \cdot \left(\frac{z}{a_0}\right)^{3/2} \int d^3x e^{i(k - P_f) \cdot x} \epsilon(k, \lambda) \cdot \nabla e^{-\frac{z}{a_0} r}$$

$$= \frac{-ie}{m \sqrt{2\omega_k \pi V^2}} \left(\frac{z}{a_0}\right)^{3/2} \int d^3x \epsilon(k, \lambda) \cdot (i(k - P_f)) e^{i(k - P_f) \cdot x} e^{-\frac{z}{a_0} r}$$

$$= \frac{e}{m\sqrt{2\pi}W_F V^2} \left(\frac{\pi}{a_0}\right)^{3/2} (\mathbf{E}(k, \lambda) \cdot \mathbf{P}_F) \int d^3x e^{i(k-P_F)x} e^{-\frac{\pi}{a_0}r}$$

$\approx \approx$

$$\int d^3x e^{i(k-P_F)x} e^{-\frac{\pi}{a_0}r} = \int_0^\infty dr r^2 \int_{-1}^1 dz \int_0^{2\pi} d\varphi e^{i|k-P_F|r z} e^{-\frac{\pi}{a_0}r}$$

$$= 2\pi \int_0^\infty dr r^2 \left[\frac{e^{i|k-P_F|r} - e^{-i|k-P_F|r}}{i|k-P_F|r} \right] e^{-\frac{\pi}{a_0}r}$$

$$= \frac{2\pi}{i|k-P_F|} \int_0^\infty dr r \left[e^{i|k-P_F|r - \frac{\pi}{a_0}r} - e^{-i|k-P_F|r - \frac{\pi}{a_0}r} \right]$$

$$= \frac{2\pi}{i|k-P_F|} \int_0^\infty dr r \left(\frac{1}{i|k-P_F| - \frac{\pi}{a_0}} \left\{ e^{i|k-P_F|r - \frac{\pi}{a_0}r} \right\}' + \frac{1}{i|k-P_F| + \frac{\pi}{a_0}} \left\{ e^{-i|k-P_F|r - \frac{\pi}{a_0}r} \right\}' \right)$$

$$= \frac{-2\pi}{i|k-P_F|} \int_0^\infty dr \left(\frac{1}{i|k-P_F| - \frac{\pi}{a_0}} e^{i|k-P_F|r - \frac{\pi}{a_0}r} + \frac{1}{i|k-P_F| + \frac{\pi}{a_0}} e^{-i|k-P_F|r - \frac{\pi}{a_0}r} \right)$$

$$= \frac{-2\pi}{i|k-P_F|} \left(\frac{1}{(i|k-P_F| - \frac{\pi}{a_0})^2} e^{i|k-P_F|r - \frac{\pi}{a_0}r} \Big|_0^\infty - \frac{1}{(i|k-P_F| + \frac{\pi}{a_0})^2} e^{-i|k-P_F|r - \frac{\pi}{a_0}r} \Big|_0^\infty \right)$$

$$= \frac{-2\pi}{i|k-P_F|} \cdot \left(-\frac{1}{(i|k-P_F| - \frac{\pi}{a_0})^2} + \frac{1}{(i|k-P_F| + \frac{\pi}{a_0})^2} \right)$$

$$= \frac{-2\pi}{i|k-P_F|} \cdot \frac{-|k-P_F|^2 - 2i\frac{\pi}{a_0}|k-P_F| + \left(\frac{\pi}{a_0}\right)^2 + |k-P_F|^2 - 2i\frac{\pi}{a_0}|k-P_F| - \left(\frac{\pi}{a_0}\right)^2}{\left(|k-P_F|^2 + \left(\frac{\pi}{a_0}\right)^2\right)^2}$$

$$= \frac{-2\pi}{i|k-P_F|} \cdot \frac{-4i\frac{\pi}{a_0}|k-P_F|}{\left(|k-P_F|^2 + \left(\frac{\pi}{a_0}\right)^2\right)^2}$$

$$= 8\pi \frac{\frac{\pi}{a_0}}{\left(|k-P_F|^2 + \left(\frac{\pi}{a_0}\right)^2\right)^2}$$

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$$\langle f | H_{int} | i \rangle = \frac{-e}{m\sqrt{2\pi}\omega_k V^2} \left(\frac{z}{a_0}\right)^{3/2} (\epsilon(k), p_f) \times 8\pi \cdot \frac{z}{a_0} \cdot \frac{1}{(|k-p_f|^2 + (\frac{z}{a_0})^2)^2}$$

$$= \frac{-8\pi e}{m\sqrt{2\pi}\omega_k V^2} \cdot \left(\frac{z}{a_0}\right)^{5/2} (\epsilon, p_f) \cdot \frac{1}{(|k-p_f|^2 + (\frac{z}{a_0})^2)^2}$$

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$$|\langle f | H_{int} | i \rangle|^2 = \frac{64\pi^2 e^2}{2\pi m^2 \omega_k V^2} \left(\frac{z}{a_0}\right)^5 (\epsilon, p_f)^2 \frac{1}{(|k-p_f|^2 + (\frac{z}{a_0})^2)^4}$$

初期状態において、電子の静止系から見ると、 $v_{rel} = 1$ (光速) といふ。

$$dV = \frac{2\pi V}{v_{rel}} \cdot V \frac{d^3 p_f}{(2\pi)^3} \delta(\omega_k - I - \frac{p_f^2}{2m}) \cdot \frac{64\pi^2 e^2}{2\pi m^2 \omega_k V^2} \left(\frac{z}{a_0}\right)^5 (\epsilon, p_f)^2 \frac{1}{(|k-p_f|^2 + (\frac{z}{a_0})^2)^4}$$

$$= d^3 p_f \delta\left(\omega_k - I - \frac{p_f^2}{2m}\right) \cdot \frac{64\pi^2 e^2}{(2\pi)^3 m^2 \omega_k} \left(\frac{z}{a_0}\right)^5 (\epsilon, p_f)^2 \frac{1}{((|k-p_f|^2 + (\frac{z}{a_0})^2)^4)}$$

$$= dP_f p_f^2 d\Omega \delta\left(\omega_k - I - \frac{p_f^2}{2m}\right) \cdot \frac{32\alpha}{m^2 \omega_k} \left(\frac{z}{a_0}\right)^5 (\epsilon, p_f)^2 \frac{1}{((|k-p_f|^2 + (\frac{z}{a_0})^2)^4)}$$

では、 $\delta(I - \omega_k - \frac{p_f^2}{2m})$ のテータ関数の処理のため、

$$dP_f = \frac{dP_f}{d(\frac{p_f^2}{2m})} d\left(\frac{p_f^2}{2m}\right) = \frac{m}{P_f} d\left(\frac{P_f^2}{2m}\right) \text{ と変形すると。} \quad \frac{d}{dp_f} \left(\frac{P_f^2}{2m}\right) = \frac{P_f}{m}$$

$$\frac{dV}{d\Omega} = \int d\Omega \delta\left(\omega_k - I - k\right) \cdot \frac{32 P_f \alpha}{m \omega_k} \left(\frac{z}{a_0}\right)^5 (\epsilon, p_f)^2 \frac{1}{((|k-p_f|^2 + (\frac{z}{a_0})^2)^4)} \quad \leftarrow \frac{P_f^2}{2m} \approx \omega_k$$

$$k \gg I, \text{無視}$$

$$= \frac{32 P_f \alpha}{m \omega_k} (\epsilon, p_f)^2 \left(\frac{z}{a_0}\right)^5 \frac{1}{\left(\left(\frac{z}{a_0}\right)^2 + (|k-p_f|^2)^2\right)^4}$$

始状態の偏極の平均をとる。

$$\begin{aligned} \frac{1}{2} \frac{\partial}{\lambda} (\delta(\mathbf{k}, \lambda) \cdot \mathbf{P}_F)^2 &= \frac{1}{2} \cdot \left(\delta_{ij} - \frac{k_i k_j}{\|\mathbf{k}\|^2} \right) P_i P_j \\ &= \frac{1}{2} \cdot \left(\mathbf{P} \cdot \mathbf{P} - \frac{(\mathbf{k} \cdot \mathbf{P}_F)^2}{\|\mathbf{k}\|^2} \right) = \\ &= \frac{1}{2} \cdot \left(\mathbf{P}_F^2 - |\mathbf{P}_F|^2 \cos^2 \theta \right) = \frac{|\mathbf{P}_F|^2}{2} (1 - \cos^2 \theta) \\ &= \frac{\mathbf{P}_F^2}{2} (1 - \cos^2 \theta) = \frac{\mathbf{P}_F^2}{2} \sin^2 \theta \end{aligned}$$

したがって、

$$\left(\frac{d\sigma}{d\Omega} \right)_{Lab} = \frac{16 \mathbf{P}_F^3 \alpha}{m \omega_k} \left(\frac{z}{a_0} \right)^5 \frac{1 - \cos^2 \theta}{\left(\left(\frac{z}{a_0} \right)^2 + (k - \mathbf{P}_F)^2 \right)^4}$$

光子の入射方向を z 軸方向に選ぶと

$$(|k - P_f|^2) = k^2 + P_f^2 - 2kP_f \cos\theta$$

$$|k| = \sqrt{k^2 + I} = \frac{P_f^2}{2m} + \frac{1}{2m} \left(\frac{z}{a_0} \right)^2 = \frac{1}{2m} \left[P_f^2 + \left(\frac{z}{a_0} \right)^2 \right]$$

\therefore

$$\left(\frac{z}{a_0} \right)^2 + (|k - P_f|^2) = \left(\frac{z}{a_0} \right)^2 + P_f^2 + k^2 - 2kP_f \cos\theta$$

$$= 2mk + k^2 - 2kP_f \cos\theta$$

$$= 2mk \left(1 + \frac{k}{2m} - \frac{P_f}{m} \cos\theta \right)$$

$$\approx 2mk (1 - v_f \cos\theta)$$

$$P_f^2 \ll 2mk \gg k^2$$

$$\underbrace{\frac{k}{2m}}_{\ll 1}$$

(a)

$$\frac{d\Gamma}{d\Omega} = \frac{16P_f^3 \alpha}{mk} \cdot \left(\frac{z}{a_0} \right)^5 \cdot \frac{1 - \cos^2\theta}{16m^4 k^4 (1 - v_f \cos\theta)^4}$$

$$= \frac{\alpha P_f^3}{5m^5 k^5} \left(\frac{z}{a_0} \right)^5 \cdot \frac{1 - \cos^2\theta}{(1 - v_f \cos\theta)^4}$$

$$= \frac{\alpha \cdot 2\sqrt{2} m^{3/2} k^{3/2}}{m^5 k^5} \left(\frac{z}{a_0} \right)^5 \cdot \frac{1 - \cos^2\theta}{(1 - v_f \cos\theta)^4}$$

$$= \frac{2\sqrt{2} \alpha}{(mk)^{1/2}} \cdot \left(\frac{z}{a_0} \right)^5 \cdot \frac{1 - \cos^2\theta}{(1 - v_f \cos\theta)^4}$$

